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# Numerical study on transport properties of the working mixtures for coal supercritical water gasification based power generation systems



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#### HIGHLIGHTS

- Viscosity of the H<sub>2</sub>O/CO<sub>2</sub>/H<sub>2</sub> mixtures and H<sub>2</sub>O/CO<sub>2</sub> mixtures is investigated.
- Predicted values by EMD method and theoretical models are compared.
- Self-diffusion coefficient, H-bond number and RDF are calculated.
- The breakdown of Stokes–Einstein relation is found and discussed.

# G R A P H I C A L A B S T R A C T



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# ABSTRACT

Viscosity of  $H_2O/CO_2/H_2$  mixtures or  $H_2O/CO_2$  mixtures is a vital transport property and required in the design of equipment for coal supercritical water gasification based power generation systems, however, no experimental data and studies have examined these viscosities in supercritical regions of water. In this paper, viscosity of these mixtures in supercritical regions of water is investigated by molecular dynamics method and different theoretical models. Moreover, the self-diffusion coefficient, radial distribution functions and H-bond number are also calculated via molecular dynamics simulations to make a better understanding of the temperature dependences of the viscosity of the mixtures in supercritical regions of water at a molecular level. The breakdown of Stokes–Einstein relation for the  $H_2O/CO_2/H_2$  mixture in supercritical regions of water is found and discussed. The prediction models and data put forth in this paper offer great value for practical application systems involving coal supercritical water gasification.

#### 1. Introduction

Clean coal technologies, such as coal gasification, have been developed around the world to minimize environmental effects from coal utilization [1–3]. In recent years, coal supercritical water gasification technology and corresponding thermodynamics cycle power generation systems have received extensive attention [4–8]. With the technique of "boiling coal in water" proposed by Guo et al., [6] the organic matter of

coal can be converted into gas (mainly  $H_2$  and  $CO_2$ ) while other elements (e.g., N, S, P, As and Hg) are deposited as inorganic salts. The two typical technology roadmaps [7,8] for the integrated thermodynamics cycle power generation systems based on coal supercritical water gasification are shown in Fig. 1. The production of the gasifier can be  $H_2O/CO_2/H_2$  or  $H_2O/CO_2$  mixtures in supercritical regions of water after separation and purification, and then they flow into the thermal power generation systems to generate electric power. This technology

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Nomenclature			Subscript			
k <sub>B</sub>	Boltzmann constant. 1.3806 $\times 10^{-23}$ J/K	i. i	refers to individual molecule $i$ or $i$			
P	system pressure, MPa	α.β	the direction in Cartesian coordinates			
V	system volume, Å <sup>3</sup>	т, <i>г</i>	refers to mixture			
Т	system temperature, K	C	critical			
t	time, s	R	reduced			
$v_i$	velocity of molecule <i>i</i> , m/s					
$m_i$	mass of molecule <i>i</i> , kg	Superscrip	bt			
r <sub>ii</sub>	distance between molecules (sites) <i>i</i> and <i>j</i> , Å	1 1				
$f_{ii}$	force acting on molecule <i>i</i> due to interactions with mole-	k, l	refers to individual site $k$ or $l$			
- 0	cule j, N	Sim	simulation data			
u <sub>ii</sub>	interaction potential energy between molecules <i>i</i> and <i>j</i> , kJ/	Exp	experimental data			
5	mol	NIST	data from national institute of standards and technology			
q	atomic charge, e					
x	mole fraction, unitless	Abbreviat	ion			
w	mass fraction, unitless					
M	molecular weight, g/mol	MD	molecular dynamics			
$T_B$	the boiling point at 1 atmosphere pressure, K	EMD	equilibrium molecular dynamics			
$N_A$	Avogadro constant, $6.02 \times 10^{23}  \text{mol}^{-1}$	RDF	the radial distribution function			
Ν	total molecular number in the system	NACF	stress-stress normalized autocorrelation function			
$C_{SE}$	the Stokes-Einstein coefficient	LJ	Lennard-Jones			
D	the self-diffusion coefficient, $m^2 s^{-1}$	LAMMPS	large-scale atomic/molecular massively parallel simulator			
$V_{vdW}$	the van-der-Waals volume, L/mol	PPPM	particle-particle/particle-mesh			
n <sub>OH</sub>	H-bond number	KRW	Kestin-Ro-Wakeham model			
R	the ideal gas constant, 8.314472 J/(K mol)	DS	Dean-Stiel model			
		EH	Ely and Hanley model			
Greek syn	nbols	NIST	national institute of standards and technology			
1.1		NEMD	non-equilibrium molecular dynamics			
$\varepsilon_{ij}^{kl}$	the LJ energy parameter, kcal/mol	ARD	absolute relative deviation			
$\sigma^{kl}_{ij}$	the LJ size parameter, Å	AARD	average absolute relative deviation			
$\varepsilon_0$	the vacuum permittivity	ARE	absolute relative error			
η	viscosity, µPa s	AARE	average absolute relative error			
ρ	density, kg/m <sup>3</sup>	MSD	the mean square displacement			

has many advantages, such as no pollutants, zero net  $CO_2$  emissions, and high coal-electricity efficiency [7].

The thermophysical properties of  $H_2O/CO_2/H_2$  or  $H_2O/CO_2$  mixtures should be understood in the design and analysis of the coal supercritical water gasification based thermodynamics cycle power generation systems [9–11]. Among the thermophysical properties, viscosity is one of the most important transport properties for the fluids. However, so far there is no report or data available for the viscosity for  $H_2O/CO_2/H_2$  mixtures and  $H_2O/CO_2$  mixtures in supercritical regions of water, and it is fairly difficult and challenging to be obtained accurately under high-temperature extreme conditions via experimental measurements [12].

Molecular dynamics (MD) simulation has been proved to be an effective theoretical way of studying the viscosity of fluids. Liang et al. [13] employed an ab initio potential model to predicted shear viscosity of CO<sub>2</sub> gas without using any experimental data at the pressure of 1 atm and temperature range of 300-1000 K; the maximum error of simulation results compared with experiment data was 1.82%. Medina et al. [14] carried out equilibrium molecular dynamics (EMD) simulations for liquid water at temperatures in the range of 273-368 K. They found that the viscosity obtained by the flexible SPC/Fw model agrees with experimental data better than that by the rigid SPC/E model. Nieszporeket al. [15] calculated the viscosity of aqueous sodium perchlorate solution. The average error of viscosity obtained using TIP4P/ 2005 water model was 9.0% compared with experiment data. Yu et al. [16] investigated the viscosity of SiO<sub>2</sub>/H<sub>2</sub>O solid-gas system using EMD simulations with three different SiO<sub>2</sub>/H<sub>2</sub>O models. In many recent studies for the thermal energy storage and concentrating solar power applications [17-19], the viscosity of the melts or the nanofluids are

also investigated via MD simulations. Ding et al. [17] performed MD simulations on viscosity of molten alkali carbonate  $K_2CO_3$ , and they found the simulation results are reasonable and reliable with an overall error of 12.05% compared with experimental data. Aguilar et al. [19] investigated the viscosity of NiO-based nanofluids with different nanoparticle mass concentrations for concentrating solar power applications. Besides the MD simulation method, some empirical theoretical models [20–24] for evaluating the viscosities of the mixtures may be another alternative, however most of these models are developed based on experimental data of common gases and their prediction reliability for water based mixtures in supercritical regions of water are still unknown.

In this paper, the viscosity of  $H_2O/CO_2/H_2$  mixtures and  $H_2O/CO_2$  mixtures in supercritical regions of water is predicted using EMD simulations and theoretical models. The radial distribution function (RDF) and the hydrogen bonds (H-bond) are calculated and analyzed for the local structure and viscosity of the water-based mixtures. The MSD and the self-diffusion coefficient are also computed via MD simulations. To allow readers to compare the obtained calculations, all numerical data and uncertainty estimates are provided in the Data in Brief.

#### 2. Methodology

#### 2.1. Green-Kubo formula

In this work, the viscosity is calculated by the Green–Kubo formulas and the EMD simulations:



Fig. 1. Schematic diagram of the integrated thermodynamics cycle power generation systems based on coal supercritical water gasification. (a) Reproduced from Ref. [8]; (b) reproduced from Ref. [7].

$$\eta = \frac{V}{3k_BT} \int_0^\infty \sum_\alpha \sum_\beta \langle P_{\alpha\beta}(0) P_{\alpha\beta}(t) \rangle dt$$
(1)

where  $\eta$  is the shear viscosity, *V* and *T* denote the volume and the temperature of the system.  $k_B$  is the Boltzmann constant equal 1.3806504 × 10<sup>-23</sup> J/K.  $\alpha$  and  $\beta$  denote the x, y and z direction in Cartesian coordinates ( $\alpha \neq \beta$ ). The angle bracket  $\langle \rangle$  denotes the average of the autocorrelation function.  $P_{\alpha\beta}(t)$  is the components of the pressure tensor of  $\alpha\beta$  direction at time *t* which can be calculated as follows [13]:

$$P_{\alpha\beta}(t) = \frac{1}{V} \left[ \sum_{i=1}^{N} m_i v_{i\alpha} v_{i\beta} + \sum_{i=1}^{N-1} \sum_{j>i}^{N} r_{ij\alpha} f_{ij\beta} \right]$$
(2)

where *N* is the number of molecules,  $m_i$  is the mass of molecule *i*.  $v_{i\alpha}$  and  $v_{i\beta}$  is velocity component of molecule *i* at  $\alpha$  and  $\beta$  direction, while *r* and *f* correspond to the displacement and force between two molecules, respectively.

#### 2.2. Potentials for molecular dynamics simulation

In this study, the combined Lennard-Jones (LJ) and Coulomb potential is adopted:

$$u_{ij} = \sum_{k=1}^{m} \sum_{l=1}^{n} \left\{ 4\varepsilon_{ij}^{kl} \left[ \left( \frac{\sigma_{ij}^{kl}}{r_{ij}^{kl}} \right)^{12} - \left( \frac{\sigma_{ij}^{kl}}{r_{ij}^{kl}} \right)^{6} \right] + \frac{q_{i}^{k} q_{j}^{l}}{4\pi\varepsilon_{0} r_{ij}^{kl}} \right\}$$
(3)

where  $u_{ij}$  is the interaction energy between molecule *i* and molecule *j*,  $\varepsilon_{ij}^{kl}$  is the LJ energy parameter,  $\sigma_{ij}^{kl}$  is the LJ size parameter,  $r_{ij}^{kl}$  is the distance between sites *k* and *l*,  $q_i^k$  is the charge on site sites *k* of molecule *i*, and  $\varepsilon_0$  is the vacuum permittivity.

To choose appropriate force field models for pure  $H_2O$ ,  $CO_2$ ,  $H_2$ , and their mixtures, 11 force fields models are examined, including SPC/E model [25], TIP4P model [26], TIP4P/2005 model [27], SPC/Fd model [28] and SPC/Fw model [28] for pure  $H_2O$ , EPM model [29], EPM2 model [29], MSM3 model [30], Cygan model [31], and TraPPE model [32] for  $CO_2$ , and the two-site model [33] for  $H_2$ . The force field parameters of the models are listed in Table 1. To describe the interactions between unlike atoms in the system, the interaction parameters of the potential are obtained via the Lorentz-Berthelot mixing rule [34] which is the most commonly used combining rule [35].

#### 2.3. Molecular simulation details

The LAMMPS software package [36] is used to perform our EMD simulations. The simulations are conducted in a three-dimensional cubic box with periodic boundary conditions in all directions. The long-

 Table 1

 Force field parameters of the models used in this work.

Molecule	Model	Site	<i>q</i> (e)	$\sigma(\text{\AA})$	$d_{\rm OM}$ (Å)	$\varepsilon$ (kcal/mol)
$H_2O$	SPC/E	0	-0.8476	3.166	_	0.1553
		Н	0.4238	_	—	_
	TIP4P	0	_	3.1536	_	0.155
		Н	0.52	_	_	_
		Μ	-1.04	_	0.15	_
	TIP4P/2005	0	_	3.1589	_	0.1852
		Н	0.5564	_	_	_
		Μ	-1.1128	_	0.1546	_
	SPC/Fw	0	-0.82	3.1654	—	0.1554
		Н	0.41	_	_	_
	SPC/Fd	0	-0.82	3.1654	_	0.1554
		Н	0.41	—	—	-
$CO_2$	EPM	С	0.6645	2.785	_	0.0576
		0	-0.33225	3.064	_	0.1651
	EPM2	С	0.6512	2.757	_	0.0559
		0	-0.3256	3.033	_	0.1602
	MSM3	С	0.594	2.785	_	0.0576
		0	-0.297	3.014	_	0.1651
	TraPPE	С	0.7	2.8	_	0.0537
		0	-0.35	3.05	—	0.1572
	Cygan	С	0.6512	2.8	—	0.05600
		0	-0.3256	3.028	—	0.15995
$H_2$	Two-site	Н	_	2.72	-	0.0199

range electrostatic interactions are computed with the particle–particle particle-mesh (PPPM) method [37] with a cutoff distance of 13 Å and an accuracy of  $10^{-4}$  in force. For LJ interactions, a cutoff distance of 13 Å is used. The time step is set as 0.4 fs for the pure fluid and the mixtures. In the simulations, the simulations initially run in NVT ensemble to equilibrate the system at a given temperature for 1 ns, then switch to the NPT ensemble if needed to control the pressure and equilibrate the system at a given pressure for another 1 ns. Finally, it is switched to the NVE ensemble, the first 1 ns are performed to relax the system, and next 2 ns or longer time are used to calculate the viscosity.

More details for the simulations can be found in Section 3, including the adopted force field models, and the simulation strategy adopted and the details of process of the EMD simulation for the pure fluids and the mixtures.

#### 2.4. Theoretical models for the calculations of viscosity of the mixtures

The Wilke model [20], DS model [21], KRW model [22,23], and EH model [24] are commonly used in the evaluation of the viscosity of the gas mixtures [38,39]. However, there is no study for these models in the prediction of the viscosity of the  $H_2O/CO_2$  binary mixtures and the ternary mixtures of  $H_2O/H_2/CO_2$  because of the absence of experimental measurements. Here the EH model cannot be used for the  $H_2O/CO_2$  binary mixtures and the ternary mixtures and the ternary mixtures of  $H_2O/H_2/CO_2$  because of the absence of the head of the value of related binary interaction parameters from the experimental data in supercritical regions of water. Thus, the left three models are studied in this study.

The Wilke model to calculate the viscosity of mixtures is as follows [20]:

$$\eta = \sum_{i=1}^{n} \frac{\eta_i}{1 + \frac{1}{x_i} \sum_{j=1, j \neq i}^{n} x_j \phi_{ij}}$$
(4)

$$\phi_{ij} = \frac{\left[1 + (\eta_i/\eta_j)^{0.5} (M_j/M_i)^{0.25}\right]^2}{2\sqrt{2} \left[1 + (M_i/M_j)\right]^{0.5}}$$
(5)

where  $\eta$  is the calculated viscosity of the mixture,  $x_i$ ,  $M_i$ , and  $\eta_i$  are mole fraction, molecular weight and viscosity of component *i*, respectively.

The equations of the DS model to calculate the viscosity of mixtures are as follows [21]:

$$(\eta - \eta^*)\xi = 10.8 \times 10^{-5} (e^{1.439\rho_R} - e^{-1.11\rho_R^{1.858}})$$
(6)

$$\eta^* \xi = 166.8 \times 10^{-5} (0.1338T_R - 0.0932)^{5/9} T_R \ge 1.5$$
(7)

$$\eta^* \xi = 34.0 \times 10^{-5} T_R^{8/9} \ T_R < 1.5$$
(8)

$$=\frac{T_{\rm cm}^{1/6}}{\left(\sum_{i=1}^{n}x_{i}M_{i}\right)^{1/2}P_{cm}^{2/3}}\tag{9}$$

ξ

In the above equations,  $T_{\rm cm}$  and  $P_{\rm cm}$  represent critical temperature and critical pressure of mixture.  $T_R$  and  $\rho_R$  are reduced temperature and reduced density.  $\eta^*$  is the viscosity of mixtures at low pressures ( < 5 bar).

$$P_{cm} = \sum_{i=1}^{n} x_i P_{ci} \tag{10}$$

$$T_R = T/T_{\rm cm} \tag{11}$$

$$T_{cm} = \sum_{i=1}^{n} x_i T_{ci} \tag{12}$$

$$\rho_R = \rho_{\rm m} / \rho_{\rm cm} \tag{13}$$

$$\rho_{\rm cm} = \frac{\sum_{i=1}^{n} x_i M_i}{\sum_{i=1}^{n} x_i V_{ci}}$$
(14)

where  $P_{\rm ci}$ ,  $T_{\rm ci}$  and  $V_{ci}$  are critical pressure, critical temperature and critical volume of component *i*, respectively.  $\rho_{\rm cm}$  and  $\rho_{\rm m}$  are critical density and density of mixtures.

The equations of the KRW model to calculate the viscosity of mixtures are as follows [22,23]:

$$\eta = \frac{\begin{vmatrix} H_{11} \cdots H_{1n} x_1 \\ H_{n1} \cdots H_{nn} x_n \\ x_1 \cdots x_n 0 \end{vmatrix}}{\begin{vmatrix} H_{11} \cdots H_{1n} \\ H_{n1} \cdots H_{nn} \end{vmatrix}}$$
(15)

$$H_{ii} = \frac{x_i^2}{\eta_i} + \sum_{\substack{k=1\\k\neq i}}^{n} \frac{2x_i x_k}{\eta_{ik}} \frac{M_i M_k}{(M_i + M_k)^2} \left(\frac{5}{3A_{ik}^*} + \frac{M_k}{M_i}\right)$$
(16)

$$H_{ij} = -\frac{2x_i x_j}{\eta_{ij}} \frac{M_i M_j}{(M_i + M_j)^2} \left(\frac{5}{3A_{ij}^*} - 1\right) \qquad i \neq j$$
(17)

In the equation, the interaction viscosity  $\eta_{ii}$  is given by

$$\eta_{ij} = \frac{5}{16} \left( \frac{k_B M_{ij} T}{\pi N_A} \right)^{\frac{1}{2}} \frac{1}{\sigma_{ij}^2 \Omega_{22}}$$
(18)

$$M_{ij} = \frac{2M_i M_j}{M_i + M_j} \tag{19}$$

where  $k_B$  and  $N_A$  are Boltzmann constant and Avogadro constant, respectively.  $A^* = \Omega_{22}/\Omega_{11}$ .  $\Omega_{11}$  and  $\Omega_{22}$  can be calculated by the formula as follows.

$$In\Omega_{22} = 0.45677 - 0.53955(InT^*) + 0.18265(InT^*)^2 - 0.03629(InT^*)^3 + 0.00241(InT^*)^4 1 < T^* < 90$$
(20)

$$In\Omega_{11} = 0.347 - 0.444(InT^*) + 0.093(InT^*)^2 - 0.010(InT^*)^3$$
  
0.5 < T\* < 25 (21)

where  $T^* = k_B T / \epsilon_{ij}$ . The energy and length scaling parameters,  $\sigma_{ij}$  and  $\epsilon_{ij}$ , are given as [40]

$$\sigma_{ij} = (1 - k_{ij,\sigma}) \sqrt{\sigma_j \sigma_i}$$
(22)

$$\varepsilon_{ij}/k_B = (1 - k_{ij,\varepsilon})\sqrt{(\varepsilon_i/k_B)(\varepsilon_j/k_B)}$$
(23)

where  $k_{ij,\sigma}$  and  $k_{ij,\varepsilon}$  are the binary interaction parameters which obtained by fitting experimental data.

# 3. Results and discussion

#### 3.1. Predicting viscosity using EMD simulation

The uncertainties of the viscosity results for different independent runs of EMD simulations are generally rather large. To solve this problem, statistics from several independent runs are necessary to obtain a reliable value result of the viscosity [41]. Besides, appropriate correlation time and enough long simulation time for the EMD simulations should be chosen carefully. In all the EMD simulations of this work, we perform ten independent simulations on pure fluid and mixtures by changing initial velocity. Finally, an average viscosity for the given temperature and pressure (or mass density) are evaluated by averaging over the values for the ten independent runs, and the error estimate is obtained by the standard error of the values for each independent run.

To choose the correlation time, the stress–stress normalized autocorrelation function (NACF) is calculated for the pure H<sub>2</sub>O, CO<sub>2</sub>, H<sub>2</sub> and their mixtures. Fig. 2 shows the NACF for the pure H<sub>2</sub>O, CO<sub>2</sub>, H<sub>2</sub> with different force field models. The NACF are calculated at temperature of 750 K and mass density of 100 kg/m<sup>3</sup> for the pure H<sub>2</sub>O and pure CO<sub>2</sub>, and temperature of 750 K and mass density of 5 kg/m<sup>3</sup> for the pure H<sub>2</sub>. The calculated NACF for different force field models for pure water is shown in Fig. 2(a). It can be observed that the NACF decays relatively fast to zero in about 5000 fs. However, for the carbon dioxide force field models, it takes long correlation times (about 8000 fs or more dependent on the force field model adopted) to sufficiently converge the NACF, as shown in Fig. 2(b). For the pure H<sub>2</sub>, correlation time about 6000 fs is needed for the convergence of NACF to zero, as shown in Fig. 2(c).

In fact, the following two points should be considered when determine the duration of the correlation time windows  $\tau$ : (1) The correlation time used to calculate the NACF should be long enough to capture its full decay; (2) longer correlation times will have larger statistical uncertainty because less data are available for its calculation. Therefore, in this work, we choose the value of  $\tau$  in simulations while NACF fully decay and almost converge to zero.

For example, with the Cygan force field model, the calculated viscosity corresponding to the correlation time for the pure CO<sub>2</sub> at temperature of 750 K and mass density of 200 kg/m<sup>3</sup> is shown in Fig. 3(a). The calculated viscosity by EMD simulations converges at 36.4 µPa s which is very close to the $\eta^{NIST}$  = 37.1 µPa s (data from NIST database). In fact, the value of  $\tau$  is also dependent on the state of the mixtures. As shown in Fig. 3(b), with the decrease of the mass density  $\rho$ , the correlation time needed increases steadily for NACF fully decay to zero and the convergence of the viscosity.

#### 3.2. System size

Here we need to examine the effect of the system size in the simulations. EMD simulations with different system sizes, the total molecular number N = 1000, 2000, 3000, 4000, and 5000, are carried out for the H<sub>2</sub>O/CO<sub>2</sub>/H<sub>2</sub> mixtures ( $w_{H_2} = 2\%$ ,  $w_{CO_2} = 40\%$ ) at pressure of 25 MPa and temperature of 700 K. For example, Fig. 4(a) shows the simulation model for system with N = 1000; the calculated viscosity corresponding to the correlation time for ten independent runs and their average is shown Fig. 4(b). To obtain the final result, we take the converged value of the averaged viscosity. The predicted results for different system sizes are shown in Fig. 5. Markers with error bars represent the average values and the corresponding standard errors for a given *N*. The solid line indicates the average over the five system sizes and the dashed lines indicate the corresponding standard error. It can be seen that the calculated viscosity using EMD is insensitive to system size. In fact, this is one of the main advantages of the EMD method for

transport properties calculation compared to the Non-equilibrium molecular dynamics (NEMD) methods [42].



Fig. 2. The NACF calculated with different force field models: (a)  $H_2O$ ; (b)  $CO_2$ ; (c)  $H_2$ .



Fig. 3. (a) Viscosity as a function of correlation time for the pure  $CO_2$  at temperature of 750 K and mass density of 200 kg/m<sup>3</sup>; (b) NACF and viscosities as a function of correlation time for the pure  $CO_2$  with different mass densities.

### 3.3. Pure fluid

Due to the absence of experimental data related to the viscosity of  $H_2O/H_2$  binary mixtures,  $H_2O/CO_2$  binary mixtures, and  $H_2O/CO_2/H_2$  ternary mixtures in supercritical regions of water, the correct selection of force fields must be obtained by comparing data from the National Institute of Standards and Technology (NIST) with simulation results of pure fluids. In fact, many studies [43–48] have taken NIST viscosity data of pure CO<sub>2</sub>,  $H_2O$ , and  $H_2$  as referents because they are relatively accurate and publicly available. In this article, viscosity data of pure  $H_2O$ ,  $CO_2$ , and  $H_2$  from the NIST database are used as referents for calculations in EMD simulations.

First, the viscosity of pure  $H_2O$  and pure  $CO_2$  is calculated using EMD simulations at a temperature of 750 K and mass density ranging from 100 to 500 kg/m<sup>3</sup> using different force field models; results are illustrated in Fig. 6a and c, respectively. The viscosity of pure  $H_2$  is calculated using EMD simulations at a temperature of 750 K and mass density ranging from 5 to 20 kg/m<sup>3</sup>, as shown in Fig. 6e. Then, to examine viscosity variations with temperature at a given pressure, the viscosity of pure  $H_2O$ , pure  $CO_2$ , and pure  $H_2$  is calculated through EMD simulations at a pressure of 25 MPa and temperature range of 700 to 950 K, as indicated in Fig. 6b, d, and f, respectively. NIST data are also presented in Fig. 6 for comparison with the simulation results.

To clarify the comparison and discussion, the absolute relative



Fig. 5. Viscosity calculated by EMD simulation for  $H_2O/H_2/CO_2$  mixtures with different system sizes.



Fig. 4. (a) Simulation system for  $H_2O/H_2/CO_2$  mixtures with total molecular number N = 1000; (b) Calculated viscosity as a function of correlation time by ten independent runs and their average.



Fig. 6. The viscosity as a function of the mass density and temperature calculated by EMD simulation with different force field models. (a) and (b): H<sub>2</sub>O; (c) and (d): CO<sub>2</sub>; (e) and (f): H<sub>2</sub>.

errors (AREs) between MD simulation results and NIST data (or experimental data) are calculated as ARE =  $|\eta^{sim} - \eta^{NIST}|/\eta^{NIST} \times 100\%$  (or ARE =  $|\eta^{sim} - \eta^{EXP}|/\eta^{EXP} \times 100\%$ ), where  $\eta^{sim}$ ,  $\eta^{NIST}$ , and  $\eta^{EXP}$  respectively denote the viscosity values from MD simulations, the NIST database and experiments. ARE results appear in Fig. 7 with corresponding data listed in Tables S2–S7 of the Data in Brief. The average

absolute relative error (AARE) of the SPC/Fw model is smallest in the force field models for pure water, with a value of 2.31%. Among the force field models for pure  $CO_2$ , the Cygan models exhibit the best accuracy with AARE value of 4.08%. The AARE of the two-site model for pure H<sub>2</sub> is 2.52%, indicating that the two-site model can accurately describe the viscosity of hydrogen gas. Given the above simulations, we



Fig. 7. AREs of viscosity calculated by EMD simulation with different force field models. (a) H<sub>2</sub>O; (b) CO<sub>2</sub>; (c) H<sub>2</sub>.

recommend the SPC/Fw model for water molecules, Cygan model for carbon dioxide molecules, and two-site model for hydrogen molecules when predicting the viscosity of their mixtures.

#### 3.4. Model validation with experimental data

No previous data are available pertaining to the viscosity of CO<sub>2</sub>/H<sub>2</sub> binary mixtures, H<sub>2</sub>O/H<sub>2</sub> binary mixtures, or H<sub>2</sub>O/CO<sub>2</sub>/H<sub>2</sub> ternary mixtures in supercritical regions of water; only experimental data [49] of CO<sub>2</sub>/H<sub>2</sub> binary mixtures at lower pressures can be found. To validate the simulation method and force field models adopted in this paper, we calculate the viscosity of CO2/H2 binary mixtures and compare the simulation results with experimental data from Mal'tsev et al. [49] A longer production time (approximately 8 ns) in the EMD simulation is needed due to the prolonged correlation time (about 6 ns) to sufficiently converge the NACF for the CO<sub>2</sub>/H<sub>2</sub> binary mixtures at lower pressures. Fig. 8 depicts the results for the viscosity of CO<sub>2</sub>/H<sub>2</sub> binary mixtures at a pressure of 0.3 MPa and respective temperatures of 500 K, 800 K, and 1100 K. The predicted values of the theoretical models are also provided for comparison. Corresponding data are listed in Tables S8 and S9 of the Data in Brief, including the AAREs of the MD simulation and theoretical models when compared with the experimental data.

The AARE of the Wilke model is smallest, with a value of 3%, and that of the MD simulations is 6.51%. Here, the MD method does not exhibit better performance than the Wilke model in predicting the viscosity of  $CO_2/H_2$  binary mixtures at relatively low pressure, but the predictive accuracy of the MD simulations is better than that of the KRW model (AARE = 10.43%) and DS model (AARE = 10.75%). This

suggests that the results of the presented EMD model are acceptable.

Additionally, in the above calculations, the predictive accuracy of the models is shown in the following order: the MS model, EMD model, KRW model, and DS model; however, such a conclusion is only obtained for  $CO_2/H_2$  mixtures at relatively low pressure. How it will be for  $CO_2/H_2$  mixtures,  $H_2O/H_2$  binary mixtures, and  $H_2O/CO_2/H_2$  ternary mixtures at temperatures and pressures in supercritical regions of water remains unknown.

#### 3.5. H<sub>2</sub>O/CO<sub>2</sub> binary mixtures

No experimental data, simulation results, or NIST data are available regarding the viscosity of H<sub>2</sub>O/CO<sub>2</sub> binary mixtures in supercritical regions. Thus, we compare our simulation results with those of theoretical models in the following states at different pressures (*P*), temperatures (*T*), mass densities ( $\rho$ ), and mass fractions: (1) *P* = 25 MPa, *T* = 700–950 K,  $w_{CO_2} = 20\%$ ; (2) *P* = 25 MPa, *T* = 700–950 K,  $w_{CO_2} = 40\%$ ; (3)  $\rho = 100 \text{ kg/m}^3$ , *T* = 700–950 K,  $w_{CO_2} = 20\%$ ; (4)  $\rho = 100 \text{ kg/m}^3$ , *T* = 700–950 K,  $w_{CO_2} = 20\%$ ; (4) results are shown in Fig. 9a–d. To clarify the comparison, the absolute relative deviations (ARDs) between the MD simulation results and calculation results using theoretical models are calculated as follows:

$$ARD = \left| \frac{\eta^{sim} - \eta^{calc}}{\eta^{sim}} \right| \times 100\%$$
(24)

where  $\eta^{sim}$  is the viscosity calculated from the MD simulations, and  $\eta^{calc}$  is the value calculated using the theoretical models. The data in Fig. 9 and the corresponding ARDs are provided in Tables S10–S13 of the Data in Brief. The averaged ARDs (AARDs) are respectively 6.93%,



Fig. 8. Viscosity as a function of the mole fraction of  $H_2(x_{H2})$  at a pressure (P) of 0.3 MPa for  $CO_2/H_2$  binary mixtures. (a) T = 500 K; (b) T = 800 K; (c) T = 1100 K.



Fig. 9. Viscosity calculated by EMD simulation and the theoretical models for  $H_2O/CO_2$  mixtures.

the KRW model is largest.

12.90%, and 21.43% for the Wilke model, KRW model, and DS model. Fig. 10 describes the viscosity variation of  $H_2O/CO_2$  mixtures with different  $w_{CO_2}$  at $\rho = 100 \text{ kg/m}^3$  and a temperature of 750 K. Data and accompanying ARDs are provided in Tables S14–S15 of the Data in Brief. According to the statistics, the overall AARDs for the Wilke model, KRW model, and DS model are 3.82%, 29.95%, and 12.77%, respectively; hence, the prediction results of the Wilke model and EMD simulations are comparable, but deviations in the KRW model and the DS model when compared with EMD simulations are relatively large.

#### 3.6. H<sub>2</sub>O/CO<sub>2</sub>/H<sub>2</sub> ternary mixtures

For H<sub>2</sub>O/CO<sub>2</sub>/H<sub>2</sub> ternary mixtures, EMD simulations are carried out and findings are compared with theoretical models in the following states: (1) P = 25 MPa, T = 700-950 K,  $w_{H_2} = 1\%$ ,  $w_{CO_2} = 20\%$ ; (2) P = 25 MPa, T = 700-950 K,  $w_{H_2} = 2\%$ ,  $w_{CO_2} = 40\%$ ; (3)  $\rho = 100$  kg/m<sup>3</sup>, T = 700-950 K,  $w_{H_2} = 1\%$ ,  $w_{CO_2} = 20\%$ ; (4)  $\rho = 100$  kg/m<sup>3</sup>, T = 700-950 K,  $w_{H_2} = 2\%$ ,  $w_{CO_2} = 40\%$ . Fig. 11a-d indicate the results; data and ARDs appear in Tables S16–S19 of the Data in Brief. The overall AARDs are respectively 3.77%, 24.72%, and 22.92% for the Wilke model, KRW model, and DS model.

Fig. 12 indicates the viscosity variation in  $H_2O/H_2/CO_2$  mixtures with different  $w_{CO_2}$  at  $\rho = 100 \text{ kg/m}^3$  and 750 K; refer to Tables S20–S21 of the Data in Brief for detailed data and ARDs. The AARDs for the Wilke model, KRW model, and DS model are 8.88%, 24.62%, and 15.27%, respectively. Overall, when predicting the viscosity of  $H_2O/H_2/CO_2$  mixtures, the AARD of the Wilke model is smallest and that of



**Fig. 10.** Viscosity as a function of  $w_{CO_2}$  calculated by EMD simulation and the theoretical models for H<sub>2</sub>O/CO<sub>2</sub> mixtures.



Fig. 11. Viscosity as a function of temperature calculated by EMD simulation and the theoretical models for H<sub>2</sub>O/H<sub>2</sub>/CO<sub>2</sub> mixtures.

# 3.7. Discussion

From the above calculations, we have provided an EMD simulation model and calculation data for the viscosity of  $H_2O/CO_2$  mixtures and  $H_2O/H_2/CO_2$  mixtures in supercritical regions of water in the absence of experimental data. Since the validity and accuracy of the adopted

MD method has been demonstrated in predicting the viscosity of pure  $H_2O$ , pure  $H_2$ , pure  $CO_2$ , and  $CO_2/H_2$  mixtures. Thus, the MD simulation results can serve as a reference to choose the theoretical models in the absence of experimental data. By comparing the EMD simulation model with the theoretical models, the AARD of the Wilke model is smallest among all models.



Fig. 12. Viscosity as a function of  $w_{CO_2}$  calculated by EMD simulation and the theoretical models for  $H_2O/H_2/CO_2$  mixtures.

For the Wilke model, component viscosity is used to predict the viscosity of the mixtures, and in our work the viscosity of pure H<sub>2</sub>O, pure H<sub>2</sub>, and pure CO<sub>2</sub> are obtained directly from the NIST database, which is close to the experimental data. In fact, Wilke [20] has stated that this model can provide satisfactory prediction on multicomponent systems if the viscosities of pure components of the mixtures are accurate at a given temperature and pressure. Moreover, Tan et al. [38] reported that the Wilke model has an AARE of 2.0% for viscosity of  $CO_2/N_2$  and  $CO_2/N_2/Ar$  mixtures at different working conditions compared to experimental data. All these indicate that the Wilke model can provide acceptable predictions.

In the KRW model, the viscosity prediction depends heavily on binary interaction parameters in the model. Unfortunately, in predicting the viscosity of H<sub>2</sub>O/CO<sub>2</sub> mixtures and H<sub>2</sub>O/H<sub>2</sub>/CO<sub>2</sub> mixtures, the binary interaction parameters  $k_{ij,\sigma}$  and  $k_{ij,e}$  for H<sub>2</sub>O/H<sub>2</sub> and H<sub>2</sub>O/ CO<sub>2</sub> are not available when using this model due to a lack of experimental data; these parameters are therefore set to zero in accordance with the suggestion in Ref [38], resulting in poor predictive accuracy. The DS model is empirical and does not require binary interaction parameters for gas pairs and the component viscosities as input. However, as pointed out by Tan et al. [38], Eqs. (7) and (8) in the DS model were developed based on data for pure argon or other nonpolar substances via regression. To achieve a more accurate prediction, a regression curve should be developed based on experimental data of the H<sub>2</sub>O/CO<sub>2</sub> and H<sub>2</sub>O/H<sub>2</sub>/CO<sub>2</sub> mixtures. Based on the above comparisons and discussion, the Wilke model can be recommended to predict the viscosity of H<sub>2</sub>O/CO<sub>2</sub> mixtures and H<sub>2</sub>O/H<sub>2</sub>/CO<sub>2</sub> mixtures among all tested theoretical models.

#### 3.8. Structure analysis

In both water and water-based mixtures, the hydrogen bonds (Hbond) affect the transportation properties [50-52]. However, It is still quite challenging to evaluate the H-bond effect on the viscosity of the water-based mixtures quantitatively, and the H-bond effect on the viscosity of the supercritical water is still unknown. Here we calculate the radial distribution function (RDF) of the H<sub>2</sub>O/H<sub>2</sub>/CO<sub>2</sub> ternary mixtures for the MD simulations in the following two states: (1) P = 25 MPa, T = 700-950 K,  $w_{H_2} = 1\%$ ,  $w_{CO_2} = 20\%$ ; (2)  $\rho = 100$  kg/ m<sup>3</sup>, T = 700-950 K,  $w_{H_2} = 1\%$ ,  $w_{CO_2} = 20\%$ , as shown in Fig. 13(a) and (b). The total RDF  $(g(r)_{total})$  is the weighted sum of g(r) for all the different atom pairs. In Fig. 13(b) and (c), the  $g(r)_{O(H_2O)-H(H_2O)}$  for this two conditions are provided to describe the O-H distances for neighboring H<sub>2</sub>O molecules, and the first peaks locates at about 2.0 Å, which are shorter than the sum of Van der Waals radii of the oxygen and hydrogen atoms of 2.7 Å. With the increase of the temperature, the amplitude of the peaks for  $g(r)_{O(H_2O)-{\rm H}(H_2O)}$  decrease gradually, and the first peaks of  $g(r)_{O(H_2O)-H(H_2O)}$  even becomes an insignificant shoulder at its place. These indicate that the size of small clusters and the average number of H-bonds per molecule decrease with increasing temperature.

H-bond can be defined by geometric conditions and energy conditions [53]: (1)  $R(O \cdot H) \leq 2.41$ Å; (2)  $\theta(O - H \cdot O) \geq 130^{\circ}$ ; (3) an interaction energy between the hydrogen bonded water molecules to be more negative than -12.9 kJ/mol (symbol  $\cdot$  indicates the H-bond between different water molecules, while symbol - indicates the Hydroxide inside the water molecule). To evaluate the H-bond effect quantitatively, the H-bond number is calculated via MD simulation. As shown in Fig. 14(a) and (b), the average H-bond number  $n_{OH}$  for both pure water and H<sub>2</sub>O/H<sub>2</sub>/CO<sub>2</sub> ternary mixtures in supercritical regions



Fig. 13. RDF between different atom types for  $H_2O/H_2/CO_2$  mixtures: (a) All atom types; (b)  $O(H_2O)-H(H_2O)$ . ( $w_{H_2} = 1\%$ ,  $w_{CO_2} = 20\%$ , P = 25 MPa).



**Fig. 14.** The average H-bond number  $n_{OH}$  as a function of temperature for both pure water and H<sub>2</sub>O/H<sub>2</sub>/CO<sub>2</sub> ternary mixtures in supercritical regions of water. (a) P = 25 MPa; (b)  $\rho = 100$  kg/m<sup>3</sup>.

decreases exponentially with the temperature. Such a trend of exponential decay of  $n_{OH}$  with the temperature for SCW agrees well with those reported data by Mizan et al. [54].

When only considering pure water in supercritical region and following the idea of Marcus [50], we can obtain the relationship between the viscosity and H-bond number for supercritical water (P = 25 MPa, T = 700-950 K; or  $\rho = 100$  kg/m<sup>3</sup>, T = 700-950 K) by regression analysis as follows:

$$\log(\eta)(\mu Pa \cdot s) = 6.02 \left(1 - \frac{V_{vdW}}{V_m}\right) - 20298.51 \frac{1}{RT} + 5.12n_{OH}$$
(25)

where  $V_{vdW}$  is the van-der-Waals volume,  $V_m$  is the molar volume, R is the ideal gas constant and equals to 8.314472 J/(K mol). The value of the  $V_{vdW}$  can be obtained from DIPPR database [55] and equals 12.4 cm<sup>3</sup>/mol. This equation can well explain the increasing trend of viscosity of the supercritical water with the increasing temperature shown in Fig. 6(b). However, it is still quite challenging to evaluate the H-bond effect on the viscosity of the water-based mixtures quantitatively because how to obtain the van-der-Waals volume  $V_{vdW}$  or the intrinsic volume of the mixtures is still questionable. Thus further investigation will be conducted on this issue in our future work.

Moreover, the Stokes–Einstein equation has been extensively used to calculate or estimate the viscosity of different fluids under ambient conditions. However, it may break down in the estimation of the viscosity for  $H_2O/H_2/CO_2$  ternary mixtures in supercritical regions.

The Stokes–Einstein equation describes the relationship between the self-diffusion coefficient and viscosity as follows:

$$D\eta/T = \frac{2k_B}{\pi C_{SE}\sigma}$$
(26)

where  $\sigma$  is the hydrodynamic diameter and  $C_{SE}$  is the Stokes-Einstein coefficient. As is well known,  $C_{SE} = 6$  for the stick boundary condition and  $C_{SE} = 4$  for the slip boundary condition.

The self-diffusion coefficient *D* can be obtained from the mean square displacement (MSD):

$$D = \frac{1}{6} \lim_{t \to \infty} \frac{d}{dt} \left\langle \frac{1}{N} \sum_{i=1}^{N} [r_i(0) - r_i(t)]^2 \right\rangle$$
(27)

where  $r_i(t)$  and  $r_i(0)$  are the position of  $i^{\text{th}}$  molecule at time *t* and *0*. The part inside the angle brackets denotes an average of MSD.

The MSD of H<sub>2</sub>O/CO<sub>2</sub>/H<sub>2</sub> mixture withw<sub>H2</sub> = 1% andw<sub>CO2</sub> = 20% is predicted with MD method for the two conditions: P = 25 MPa, T = 700-950 K; (2)  $\rho = 100$  kg/m<sup>3</sup>, T = 700-950 K, as shown in Fig. 15(a) and (c). By using Eq. (27), the self-diffusion coefficient *D* is calculated from the MSD values, as shown in Fig. 15(b) and (d). It can be observed that the self-diffusion coefficient *D* almost increase lineally with the increasing temperature.

The Stokes–Einstein relation states that  $D\eta/T$  is a constant and is independent of the external conditions. However, studies [56,57] have shown the breakdown of Stokes–Einstein relation in the supercooled liquids or supercritical water. Here we calculated the  $D\eta/T$  for the H<sub>2</sub>O/CO<sub>2</sub>/H<sub>2</sub> mixture ( $w_{H_2} = 1\%$  and $w_{CO_2} = 20\%$ ), as shown in Fig. 16. We find that it doesn't obey the Stokes–Einstein relation. In H<sub>2</sub>O/CO<sub>2</sub>/H<sub>2</sub> mixtures at the studied conditions,  $D\eta/T$  increases lineally with the increasing of the temperature. Therefore we can conclude that it is not suitable to predict the viscosity of H<sub>2</sub>O/CO<sub>2</sub>/H<sub>2</sub> mixtures or H<sub>2</sub>O/CO<sub>2</sub> mixtures by directly using the Stokes–Einstein relation.

#### 4. Conclusions

The viscosity of  $H_2O/CO_2$  and  $H_2O/CO_2/H_2$  mixtures is one of the key transport properties in design and optimization of the equipments in thermodynamics cycle power generation system based on coal supercritical water gasification. In this paper, the viscosity properties of the  $H_2O/CO_2$  binary mixture and  $H_2O/CO_2/H_2$  mixtures in supercritical regions of water are investigated via EMD simulations and theoretical methods. Simulation models for predicting viscosity of the mixtures are validated, and force models and simulation strategy are recommended. Comparisons show that the prediction results of the Wilke model agree well with the MD simulations.

The RDF and the hydrogen bonds are calculated and analyzed for the local structure and viscosity of the water-based mixtures. The MSD and the self-diffusion coefficient are also computed via MD simulations, and the breakdown of the Stokes–Einstein relation in  $H_2O/CO_2/H_2$ mixtures in supercritical regions of water is found and discussed. Results in this paper could offer references for the design and optimization of a thermodynamic system based on coal supercritical water gasification.

#### **Declaration of Competing Interest**

The authors declare no competing financial interest.



Fig. 15. The MSD and the self-diffusion coefficient Dof  $H_2O/CO_2/H_2$  mixtures( $w_{H_2} = 1\%$  and  $w_{CO_2} = 20\%$ ).



Fig. 16. The calculated  $D\eta/T$  for the H<sub>2</sub>O/CO<sub>2</sub>/H<sub>2</sub> mixture ( $w_{H_2} = 1\%$  and  $w_{CO_2} = 20\%$ ).

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